An Explicit Construction of Quantum Expanders

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Abstract

Quantum expanders are a natural generalization of classical expanders. These objects were introduced and studied by [1, 3, 4]. In this note we show how to construct explicit, constant-degree quantum expanders. The construction is essentially the classical Zig-Zag expander construction of [5], applied to quantum expanders.

1 Introduction

Classical expanders are graphs of low degree and high connectivity. One way to measure the expansion of a graph is through the second eigenvalue of its adjacency matrix. This paper investigates the quantum counterpart of these objects, defined as follows. For a linear space $\mathcal{V}$ we denote by $L(\mathcal{V})$ the space of linear operators from $\mathcal{V}$ to itself.

Definition 1.1. *We say an admissible superoperator $G : L(\mathcal{V}) \to L(\mathcal{V})$ is D-*regular *if $G = \frac{1}{D}\sum_d G_d$, and for each $d \in [D]$, $G_d(X) = U_d X U_d^\dagger$ for some unitary transformation U_d over $\mathcal{V}$.*

Definition 1.2. *An admissible superoperator $G : L(\mathcal{V}) \to L(\mathcal{V})$ is a $(N, D, \overline{\lambda})$ quantum expander if $\dim(\mathcal{V}) = N$, G is D-regular and:*

- *$G(\tilde{I}) = \tilde{I}$, where $\tilde{I}$ denotes the completely-mixed state.*
- *For any $\rho \in L(\mathcal{V})$ that is orthogonal to $\tilde{I}$ (with respect to the Hilbert-Schmidt inner product, i.e. $\mathrm{Tr}(\rho\tilde{I}) = 0$) it holds that $\| G(A) \| \le \overline{\lambda} \| A \|$ (where $\| X \| = \sqrt{\mathrm{Tr}(XX^\dagger)}$).*

A quantum expander is explicit *if G can be implemented by a quantum circuit of size polynomial in $\log(N)$.*

The notion of quantum expanders was introduced and studied by [1, 3, 4]. These papers gave several constructions and applications of these objects. The disadvantage of all the constructions given by these papers is that each construction is either constant-degree or explicit, but not both. In this paper we show how to construct explicit quantum expanders of constant-degree. Our construction is an easy generalization of the Zig-Zag expander construction given in [5].

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2 Preliminaries

We denote by $\mathcal{H}_N$ the Hilbert space of dimension N.

For a linear space $\mathcal{V}$, we denote by $L(\mathcal{V})$ the space of linear operators from $\mathcal{V}$ to itself. We use the Hilbert-Schmidt inner product on this space, i.e. for $X, Y \in L(\mathcal{V})$ their inner product is $\langle X, Y \rangle = \mathrm{Tr}(XY^\dagger)$. The inner product gives rise to a norm $\| X \| = \sqrt{\langle X, X \rangle} = \sqrt{\sum s_i(X)^2}$, where $\{s_i(X)\}$ are the singular values of X. Throughout the paper this is the only norm we use.

We also denote by $U(\mathcal{V})$ the set of all unitary operators on $\mathcal{V}$, and by $T(\mathcal{V})$ the space of superopeartors on $\mathcal{V}$ (i.e. $T(\mathcal{V}) = L(L(\mathcal{V}))$).

Finally, we denote by $\tilde{I}$ the identity operator normalized such that $\mathrm{Tr}(\tilde{I}) = 1$. That is, $\tilde{I}$ denotes the completely mixed state (on the appropriate space).

3 Explicit constant-degree quantum expanders

3.1 The basic operations

The construction uses as building blocks the following operations:

- **Squaring:** For a superoperator $G \in T(\mathcal{V})$ we denote by G^2 the superoperator given by $G^2(X) = G(G(X))$ for any $X \in L(\mathcal{V})$.
- **Tensoring:** For superoperators $G_1 \in T(\mathcal{V}_1)$ and $G_2 \in T(\mathcal{V}_2)$ we denote by $G_1 \otimes G_2$ the superoperator given by $(G_1 \otimes G_2)(X \otimes Y) = G_1(X) \otimes G_2(Y)$ for any $X \in L(\mathcal{V}_1), Y \in L(\mathcal{V}_2)$.
- **Zig-Zag product:** For superoperators $G_1 \in T(\mathcal{V}_1)$ and $G_2 \in T(\mathcal{V}_2)$ we denote by $G_1 \textcircled{z} G_2$ their Zig-Zag product. A formal definition of this is given in Section 4. The only requirement is that G_1 is $\dim(\mathcal{V}_2)$-regular.

Proposition 3.1. *If G is a (N, D, λ) quantum expander then G^2 is a (N, D^2, λ^2) quantum expander. If G is explicit then so is G^2.*

Proposition 3.2. *If G_1 is a (N_1, D_1, λ_1) quantum expander and G_2 is a (N_2, D_2, λ_2) quantum expander then $G_1 \otimes G_2$ is a $(N_1 \cdot N_2, D_1 \cdot D_2, \max(\lambda_1, \lambda_2))$ quantum expander. If G_1 and G_2 are explicit then so is $G_1 \otimes G_2$.*

Theorem 1. *If G_1 is a (N_1, D_1, λ_1) quantum expander and G_2 is a (D_1, D_2, λ_2) quantum expander then $G_1 \textcircled{z} G_2$ is a $(N_1 \cdot D_1, D_2^2, \lambda_1 + \lambda_2 + \lambda_2^2)$ quantum expander. If G_1 and G_2 are explicit then so is $G_1 \textcircled{z} G_2$.*

The proofs of Propositions 3.1 and 3.2 are trivial. The proof of Theorem 1 is given in Section 4.

3.2 The construction

The construction starts with some constant-degree quantum expander, and iteratively increases its size via alternating operations of squaring, tensoring and Zig-Zag products. The tensoring is used to square the dimension of the superoperator. Then a squaring operation improves the second eigenvalue. Finally, the Zig-Zag product reduces the degree, without deteriorating the second eigenvalue too much.

Suppose H is a (D^8, D, λ) quantum expander. We define a series of superoperators as follows. The first two superoperators are $G_1 = H^2$ and $G_2 = H \otimes H$. For every $t > 2$ we define

$$G_t = \left(G_{\lceil \frac{t-1}{2} \rceil} \otimes G_{\lfloor \frac{t-1}{2} \rfloor} \right)^2 ⓩ H.$$

Theorem 2. *For every $t > 0$, G_t is an explicit (D^{8t}, D^2, λ_t) quantum expander with $\lambda_t = \lambda + O(\lambda^2)$.*

The proof of this Theorem for classical expanders was given in [5]. The proof only relies on the properties of the basic operations. Proposition 3.1, Proposition 3.2 and Theorem 1 assure the required properties of the basic operations are satisfied in the quantum case as well. Hence, the proof of this theorem is identical to the one in [5] (Theorem 3.3) and we omit it.

3.3 The base superoperator

Theorem 2 relies on the existence of a good base superoperator H. In the classical setting, the probabilistic method assures us that a good base graph exists, and so we can use an exhaustive search to find one. The quantum setting exhibits a similar phenomena:

Theorem 3. *([4]) There exists a D_0 such that for every $D > D_0$ there exist a (D^8, D, λ) quantum expander for $\lambda = \frac{4\sqrt{D-1}}{D}$* [1]*.*

We will use an exhaustive search to find such a quantum expander. To do this we first need to transform the searched domain from a continuous space to a discrete one. We do this by using a net of unitary matrices, $S \subset U(\mathcal{H}_{D^8})$. S has the property that for any unitary matrix $U \in U(\mathcal{H}_{D^8})$ there exists some $V_U \in S$ such that

$$\sup_{\| X \|=1} \left\| UXU^\dagger - V_U X V_U^\dagger \right\| \leq \lambda.$$

It is not hard to verify that indeed such S exists, with size depending only on D and λ. Moreover, we can find such a set in time depending only on D and λ [2].

Suppose G is a (D^8, D, λ) quantum expander, $G(X) = \frac{1}{D}\sum_{i=1}^{D} U_i X U_i^\dagger$. We denote by G' the superoperator $G'(X) = \frac{1}{D}\sum_{i=1}^{D} V_{U_i} X V_{U_i}^\dagger$. Let $X \in L(\mathcal{H}_{D^8})$ be orthogonal to $\tilde{I}$. Then:

$$\begin{aligned} \| G'(X) \| = \left\| \frac{1}{D}\sum_{i=1}^{D} V_{U_i} X V_{U_i}^\dagger \right\| &\leq \left\| \frac{1}{D}\sum_{i=1}^{D} U_i X U_i^\dagger \right\| + \frac{1}{D}\sum_{i=1}^{D} \left\| U_i X U_i^\dagger - V_{U_i} X V_{U_i}^\dagger \right\| \\ &\leq \| G(X) \| + \lambda \| X \| \leq 2\lambda \| X \|. \end{aligned}$$

Hence, G' is a $(D^8, D, \frac{8\sqrt{D-1}}{D})$ quantum expander [3]. This implies that we can find a good base superoperator in time which depends only on D and λ.

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[1] [4] actually shows that for any D there exist a $(D^8, D, (1 + O(D^{-16/15} \log D))\frac{2\sqrt{D-1}}{D})$ quantum expander.

[2] One way to see this is using the Solovay-Kitaev theorem (see, e.g., [2]). The theorem assures us that, for example, the set of all the quantum circuits of length $O(\log^4 \epsilon^{-1})$ generated only by Hadamard and Tofolli gates give an ϵ-net of unitaries. The accuracy of the net is measured differently in the Solovay-Kitaev theorem, but it can be verified that the accuracy measure we use here is roughly equivalent.

[3] We can actually get an eigenvalue bound of $(1+\epsilon)\frac{2\sqrt{D-1}}{D}$ for an arbitrary small ϵ on the expense of increasing D_0.

4 The Zig-Zag product

Suppose G_1, G_2 are two superoperators, $G_i \in T(\mathcal{H}_{N_i})$, and G_i is a (N_i, D_i, λ_i) quantum expander. We further assume that $N_2 = D_1$. G_1 is D_1–regular and so it can be expressed as $G_1(X) = \frac{1}{D_1}\sum_d U_d X U_d^\dagger$ for some unitaries $U_d \in U(\mathcal{H}_{N_1})$. We lift the ensemble $\{U_d\}$ to a superoperator $\dot{U} \in L(\mathcal{H}_{N_1} \otimes \mathcal{H}_{D_1})$ defined by:

$$\dot{U}(|a\rangle \otimes |b\rangle) \quad = \quad U_b\,|a\rangle \otimes |b\rangle\,,$$

and we define $\dot{G_1} \in T(\mathcal{H}_{N_1} \otimes \mathcal{H}_{D_1})$ by $\dot{G_1}(X) = \dot{U}X\dot{U}^\dagger$.

Definition 4.1. *Let G_1, G_2 be as above. The Zig-Zag product, $G_1 ⓩ G_2 \in T(\mathcal{H}_{N_1} \otimes \mathcal{H}_{D_1})$ is defined to be $(G_1 ⓩ G_2)X = (I \otimes G_2)\dot{G_1}(I \otimes G_2^\dagger)X$.*

We claim:

Proposition 4.2. *For any $X, Y \in L(\mathcal{H}_{N_1} \otimes \mathcal{H}_{D_1})$ such that X is orthogonal to the identity operator we have:*

$$|\ \langle\ G_1 ⓩ G_2 X\ ,\ Y\rangle\ | \quad \leq \quad f(\lambda_1, \lambda_2)\,\|\,X\,\| \cdot \|\,Y\,\|$$

where $f(\lambda_1, \lambda_2) = \lambda_1 + \lambda_2 + \lambda_2^2$.

And as a direct corollary we get:

Theorem 1. *If G_1 is a (N_1, D_1, λ_1) quantum expander and G_2 is a (D_1, D_2, λ_2) quantum expander then $G_1 ⓩ G_2$ is a $(N_1 \cdot D_1, D_2^2, \lambda_1 + \lambda_2 + \lambda_2^2)$ quantum expander. If G_1 and G_2 are explicit then so is $G_1 ⓩ G_2$.*

Proof: Let X be orthogonal to $\tilde{I}$ and let $Y = (G_1 ⓩ G_2)X$. By Proposition 4.2 $\|\,Y\,\|^2 \leq f(\lambda_1, \lambda_2)\,\|\,X\,\| \cdot \|\,Y\,\|$. Equivalently, $\|\,(G_1 ⓩ G_2)X\,\| \leq f(\lambda_1, \lambda_2)\,\|\,X\,\|$ as required.

The explicitness of $G_1 ⓩ G_2$ is immediate from the definition of the Zig-Zag product. ∎

We now turn to the proof of Proposition 4.2. We adapt the proof given in [5] for the classical case to the quantum setting. For that we need to work with linear operators instead of working with vectors. Consequently, we replace the vector inner-product used in the classical proof with the Hilbert-Schmidt inner product on linear operators, and replace the Euclidean norm on vectors, with the $\mathrm{Tr}(XX^\dagger)$ norm on linear operators. Interestingly, the same proof carries over to this generalized setting. One can get the proof below by simply going over the proof in [5] and doing the above translation. We provide the details here for completeness.

Proof of Proposition 4.2: We first decompose the space $L(\mathcal{H}_{N_1} \otimes \mathcal{H}_{D_1})$ to

$$\begin{aligned} W^{||} &= \mathrm{Span}\left\{\sigma \otimes \widetilde{I} \mid \sigma \in L(\mathcal{H}_{N_1})\right\} \text{ and,} \\ W^{\perp} &= \mathrm{Span}\left\{\sigma \otimes \tau \mid \sigma \in L(\mathcal{H}_{N_1})\ ,\ \tau \in L(\mathcal{H}_{D_1})\ ,\ \left\langle \tau, \widetilde{I}\right\rangle = 0\right\}. \end{aligned}$$

Decompose X to $X = X^{||} + X^{\perp}$, where $X^{||} \in W^{||}$ and $X^{\perp} \in W^{\perp}$, and similarly $Y = Y^{||} + Y^{\perp}$. By definition,

$$|\langle G_1 ⓩ G_2 X, Y\rangle| = \left|\left\langle (I\otimes G_2)\dot{G}_1(I\otimes G_2^{\dagger})X, Y\right\rangle\right| = \left|\left\langle \dot{G}_1(I\otimes G_2)(X^{||}+X^{\perp}), (I\otimes G_2)(Y^{||}+Y^{\perp})\right\rangle\right|.$$

Opening to the four terms and pushing the absolute value inside, we see that

$$\begin{aligned}
|\langle G_1 ⓩ G_2 X, Y\rangle| \;&\leq\; \left|\left\langle \dot{G}_1(I\otimes G_2)X^{||}, (I\otimes G_2)Y^{||}\right\rangle\right| + \left|\left\langle \dot{G}_1(I\otimes G_2)X^{||}, (I\otimes G_2)Y^{\perp}\right\rangle\right| + \\
&\quad\; \left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, (I\otimes G_2)Y^{||}\right\rangle\right| + \left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, (I\otimes G_2)Y^{\perp}\right\rangle\right| \\
&=\; \left|\left\langle \dot{G}_1X^{||}, Y^{||}\right\rangle\right| + \left|\left\langle \dot{G}_1X^{||}, (I\otimes G_2)Y^{\perp}\right\rangle\right| + \\
&\quad\; \left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, Y^{||}\right\rangle\right| + \left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, (I\otimes G_2)Y^{\perp}\right\rangle\right|
\end{aligned}$$

Where the last equality is due to the fact that $I \otimes G_2$ is identity over $W^{||}$ (since $G_2(\tilde{I}) = \tilde{I}$). In the last three terms we have $I \otimes G_2$ acting on an operator from $W^{\perp}$. As expected, when this happen the quantum expander G_2 shrinks the operator. Formally,

Claim 4.3. *For any $Z \in W^{\perp}$ we have $\|(I \otimes G_2)Z\| \leq \lambda_2 \|Z\|$.*

We defer the proof for later. Having the claim we see that, e.g., $\left|\left\langle \dot{G}_1X^{||}, (I\otimes G_2)Y^{\perp}\right\rangle\right| \leq \left\|\dot{G}_1X^{||}\right\| \cdot \|(I\otimes G_2)Y^{\perp}\| \leq \lambda_2 \|X^{||}\| \cdot \|Y^{\perp}\|$. Similarly, $\left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, Y^{||}\right\rangle\right| \leq \lambda_2 \|X^{\perp}\| \cdot \|Y^{||}\|$ and $\left|\left\langle \dot{G}_1(I\otimes G_2)X^{\perp}, (I\otimes G_2)Y^{\perp}\right\rangle\right| \leq \lambda_2^2 \|X^{\perp}\| \, \|Y^{\perp}\|$.

To bound the first term, we notice that on inputs from $W^{||}$ the operator $\dot{G}_1$ mimics the operation of G_1 with a random seed. Formally,

Claim 4.4. *For any $A \in W^{||}$ orthogonal to the identity operator and any $B \in W^{||}$ we have $\left|\left\langle \dot{G}_1A, B\right\rangle\right| \leq \lambda_1 \|A\| \cdot \|B\|$.*

We again defer the proof for later. Having the claim we see that $\left|\left\langle \dot{G}_1X^{||}, Y^{||}\right\rangle\right| \leq \lambda_1 \left\|X^{||}\right\| \cdot \left\|Y^{||}\right\|$. Denoting $p_i = \frac{\left\|\rho_i^{||}\right\|}{\|\rho_i\|}$ and $q_i = \frac{\|\rho_i^{\perp}\|}{\|\rho_i\|}$ (for $i = 1, 2$, $\rho_1 = X$ and $\rho_2 = Y$) we see that $p_i^2 + q_i^2 = 1$, and,

$$|\langle (G_1 ⓩ G_2)X, Y\rangle| \;\leq\; (p_1p_2\lambda_1 + p_1q_2\lambda_2 + p_2q_1\lambda_2 + q_1q_2\lambda_2^2)\,\|X\| \cdot \|Y\|$$

Elementary calculus now shows that this is bounded by $f(\lambda_1, \lambda_2)\,\|X\| \cdot \|Y\|$. ∎

We still have to prove the two claims:

Proof of Claim 4.3: Z can be written as $Z = \sum_i \sigma_i \otimes \tau_i$, where each τ_i is perpendicular to $\tilde{I}$ and $\{\sigma_i\}$ is an orthogonal set. Hence,

$$\|(I\otimes G_2)Z\| = \left\|\sum_i \sigma_i \otimes G_2(\tau_i)\right\| \leq \sum_i \|\sigma_i \otimes G_2(\tau_i)\| \leq \sum_i \lambda_2 \|\sigma_i \otimes \tau_i\| = \lambda_2 \|Z\|.$$

∎

And,

Proof of Claim 4.4: Since $A, B \in W^{||}$, they can be written as

$$A = \sigma \otimes \tilde{I} = \frac{1}{D_1} \sum_i \sigma \otimes |i\rangle\langle i|$$

$$B = \eta \otimes \tilde{I} = \frac{1}{D_1} \sum_i \eta \otimes |i\rangle\langle i| \,.$$

Moreover, since A is perpendicular to the identity operator, it follows that σ is perpendicular to the identity operator on the space $L(\mathcal{H}_{N_1})$. This means that applying G_1 on σ will shrink it by at least a factor of λ_1.

Considering the inner product

$$\begin{aligned}
| \left\langle \dot{G}_1 A, B \right\rangle | &= \frac{1}{D_1^2} \left| \sum_{i,j} \mathrm{Tr} \left(\left((U_i \sigma U_i^\dagger) \otimes |i\rangle\langle i| \right) (\eta \otimes |j\rangle\langle j|)^\dagger \right) \right| \\
&= \frac{1}{D_1^2} \left| \sum_{i,j} \mathrm{Tr} \left((U_i \sigma U_i^\dagger \eta^\dagger) \otimes |i\rangle \langle i\,|\,j\rangle \langle j| \right) \right| \\
&= \frac{1}{D_1^2} \left| \sum_{i} \mathrm{Tr} \left((U_i \sigma U_i^\dagger \eta^\dagger) \otimes |i\rangle\langle i| \right) \right| \\
&= \frac{1}{D_1^2} \left| \sum_{i} \mathrm{Tr} \left(U_i \sigma U_i^\dagger \eta^\dagger \right) \right| \\
&= \frac{1}{D_1} \left| \mathrm{Tr} \left(\left(\frac{1}{D_1} \sum_i U_i \sigma U_i^\dagger \right) \eta^\dagger \right) \right| \\
&= \frac{1}{D_1} \left| \langle G_1(\sigma), \eta \rangle \right| \le \frac{\lambda_1}{D_1} \parallel \sigma \parallel \cdot \parallel \eta \parallel = \lambda_1 \parallel A \parallel \cdot \parallel B \parallel ,
\end{aligned}$$

where the inequality follows from the expansion property of G_1 (and Cauchy-Schwartz). ∎

References

[1] A. Ben-Aroya and A. Ta-Shma. Quantum expanders and the quantum entropy difference problem. arXiv:quant-ph/0702129.

[2] C. M. Dawson and M. A. Nielsen. The Solovay-Kitaev algorithm. *Quantum Information & Computation*, 6(1):81–95, 2006.

[3] M. B. Hastings. Entropy and entanglement in quantum ground states. *Phys. Rev. B*, 76(3):035114, 2007.

[4] M. B. Hastings. Random unitaries give quantum expanders. Technical report, arXiv:0706.0556, 2007.

[5] O. Reingold, S. Vadhan, and A. Wigderson. Entropy waves, the zig-zag graph product, and new constant-degree expanders. *Ann. of Math.*, 155(1):157–187, 2002.